DAMPING INDUCED BY DRY FRICTION: ANALYSES AND EXPERIMENTS FOR MODELING IMPROVEMENT

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Abstract. The present study deals with the effects of contact characteristics on the vibratory response of a beam associated with several rubbing devices. Experimentally different contact configurations are tested by varying the contact area size and roughness of samples. For each tests, the frequency response functions (FRF) are measured and compared. Numerically, temporal calculations are presented and compared to measured data, using Masing's contact model with updated parameter values and averaged estimated values.

1 INTRODUCTION

The optimization of the vibratory response of structures in presence of rubbing is always very delicate due to the highly nonlinear behavior induced by dry friction occurring at the interfaces. Sticking, macro or micro slipping between interfaces have major effects on both natural frequencies and damping properties.

Due to the diversity of contact configurations, friction models are often very case dependent and cannot be efficiently used without identifying: the type of friction occurring at the interfaces (based on the shape of the hysteresis curve) and contact parameters such as friction coefficient or joint stiffness. Those information may be obtained from specific experimental analyses or from available databases. In [1] for example, the authors propose estimated friction coefficients and contact stiffness values based on a very large experimental campaign they carried.

There is always a need for better understanding of rubbing mechanisms leading to optimal configurations and for more physics-based models than the classical and simple Coulomb or Masing models. A promising way concerns models that take into account material and geometrical properties of contacting surfaces, surface roughness data, and contact loading [2-4]. These models are based on the pioneer work of Greenwood and Williamson on rough contacts [5]. Knowing the behavior of an individual spherical asperity pressed against a flat surface it is possible, thanks to statistical summation, to simulate the global behavior of two rough flat-on-flat rubbing surfaces.

In the present study, a test rig has been developed to measure the dynamic response of a steel beam associated with several friction configurations. The set-up allows punctual or surface to surface contact cases. In the case of surface to surface contact, devices with various area size, shape and roughness were tested in order to assess their influence on the overall response. Numerically, temporal calculations are presented and compared to the measured data using different contact models with updated parameter values and averaged estimated values.

2 THE TEST RIG

The studied structure is a cantilever steel beam with a rectangular cross section of 0.008 m by 0.04 m and a length of 0.50 m. At the free end, the beam may be associated with different rubbing devices (Figure 1).

The experimental set-up allows the measurement of the displacement at 0.35 m of the clamped end of the beam by mean of an accelerometer (measured signal integrated twice). The excitation force provided by a shaker is measured by a force sensor located at the end of the shaker's push rod. The normal and tangential contact forces are also given by a force sensor. Two categories of devices were chosen in order to allow punctual contacts or surface to surface contacts (Figure 2). Sine sweeps (harmonic excitation) were performed at constant excitation force amplitude within the [15-130] Hz frequency range and for various normal contact forces.

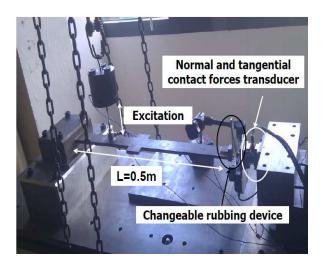


Figure 1: Experimental set up – cantilever steel beam.

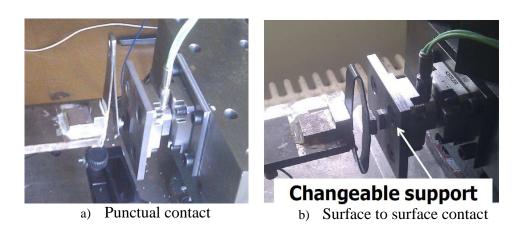


Figure 2: Rubbing devices.

Figure 3 gives an example of typical measured FRFs for the punctual contact case. The classical influence of normal contact load is exhibited.

For a high normal load (Fn=40N associated to an excitation force of 2 N) contact is stuck. The FRF exhibits a linear behaviour with a peak resonance at 86 Hz. As the normal load is lowered (from 30 N to 0.25 N) relative displacement is allowed at the contact (presence of stick/slip) and vibrational amplitudes around the 86 Hz resonance peak are damped, becoming nearly flat. Finally, for Fn=0.1 N contact is in a slip state and a high resonance peak is again observed at 24Hz (free configuration).

In Figure 4, an example of the measured tangential forces in stick and stick/slip states is shown. When contact is stuck, the tangential force is proportional to displacement (sinusoidal signal). On the other hand, when contact is in a stick/slip state, the tangential force is proportional to displacement in the stuck phase and keep the limit value of $\pm \mu Fn$ (μ friction coefficient) in the slip phase.

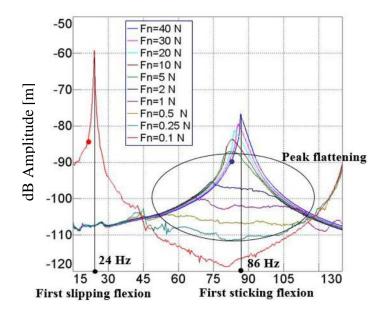


Figure 3: FRF vs normal load.

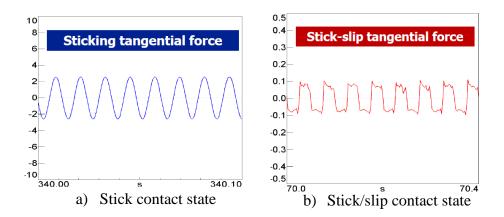


Figure 4: Measured tangential force.

3 DYNAMIC RESPONSE OF THE STRUCTURE

Only the results associated to surface to surface contact configurations are presented here. The different samples considered were made of various contact areas (13x13 mm², 9x9 mm², 13x5 mm², 5x5 mm², Figure 5) and two surface finish, one rough and one smooth. Surface finish is described by the average roughness (Ra) defined in [6]:

$$R_a = \frac{1}{L} \int_0^L |z(x)| dx \tag{1}$$

where z(x) represents the height of asperities and L the sampling length.

In Figure 6, the profiles of the $13x13 \text{ mm}^2$ samples are presented. The smooth one has a Ra of $0.714 \mu m$ while the rough one has a Ra of $3.76 \mu m$. Regarding the external forces, nor-

mal loads were varied within the range [0.1 - 80] N and the excitation force amplitude remained at 2N.

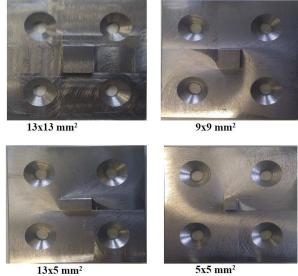


Figure 5: Samples

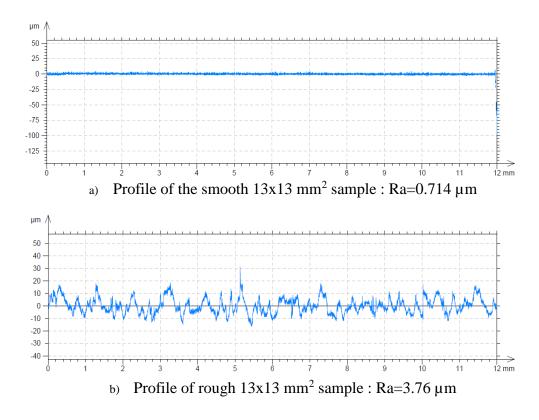


Figure 6: Profiles of the 13x13 mm² specimens.

Figure 7 gives the FRFs associated to the four smooth samples for a normal load of 2N. As clearly shown responses are identical and contact area size does not affect significantly the dynamics of the assembly. The same observations were made when assessing the influence of surface roughness.

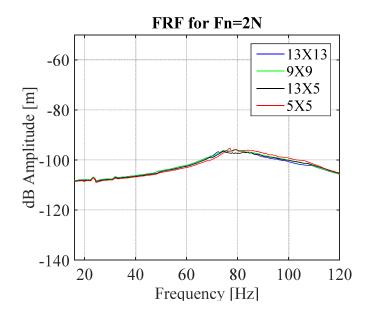


Figure 7: Area size influence on the FRFs for the smooth samples with Fn=2 N.

The overall effect of area size is illustrated in Tables 1 and 2 while the effect of roughness is shown in Tables 3 and 4. Results reported concern surfaces 13x13 mm² and 5x5 mm². Discrepancies are illustrated via three quantities: maximum displacement amplitudes, frequency at which the maximum amplitudes are obtained and damping ratios.

	Fn=0.1 N	Fn=2 N	Fn=10 N	Fn=80 N
Amplitude differ-	-5.2 %	0.7 %	1.5 %	0.8 %
ences.	3.2 70	0.7 70	1.5 70	0.0 70
Frequency at max amplitudes.	0 %	3 %	-1.1 %	0 %
Damping ratio differences.		-0.44 %	0.91 %	0.27 %

Table 1: Effect of area size: 13x13 mm² smooth versus 5x5 mm² smooth.

-	Fn=0.1 N	Fn=2 N	Fn=10 N	Fn=80 N
Amplitude differ-	4 %	1.7 %	1.3 %	3.3 %
ences.	. , ,	21,70	110 70	0.0 /0
Frequency at max amplitudes.	0 %	0 %	0 %	1.6%
Damping ratio differences.		6.9 %	0.59 %	0.81 %

Table 2: Effect of area size: 13x13 mm² rough versus 5x5 mm² rough.

As shown in Tables 1 and 2, discrepancies between surfaces 13x13 and 5x5 are limited. Variation in maximum displacement amplitudes does not exceed 5.2 % and frequencies of the

maximum amplitudes as well as damping are not affected significantly, with one exception for Fn=2 N with 6.9 % discrepancy on damping.

	Fn=0.1 N	Fn=2 N	Fn=10 N	Fn=80 N
Amplitude differences.	-0.31 %	2.4 %	0.85 %	-0.2 %
Frequency of the maximum amplitudes.	0 %	5.8 %	2.8 %	1.6 %
Damping ratio differences.		-11.3 %	-1 %	0 %

Table 3: Influence of roughness: 13x13 mm² smooth versus 13x13 mm² rough.

	Fn=0.1 N	Fn=2 N	Fn=10 N	Fn=80 N
Amplitude differences.	8.4 %	-1.5 %	-1 %	2.32 %
Frequency of the maximum amplitudes.	0 %	5.8 %	4 %	1.6 %
Damping ratio differences.		-4.1 %	-1.36 %	0.81 %

Table 4: Influence of roughness: 5x5 mm² smooth versus 5x5 mm² rough.

The former conclusions remain when considering the effects of roughness on the dynamics of the assembly as shown in Tables 3 and 4.

4 NUMERICAL MODELLING

A finite element model of the structure has been constructed using beam, shell and 8 nodes brick elements. For the contact interface, zero thickness elements are used with Masing's tangential contact model. Masing's model is a combination of a Coulomb's contact model with a spring representing contact stiffness as illustrated in Figure 8 for a single dof mass-spring system. The tangential contact force Ft is given by:

$$F_t = sgn(\dot{x})\mu Fn$$
 for slipping contact
 $F_t = -sgn(\dot{x})\mu Fn + k_d(x-z)$ for sticking contact (2)

where μ is the friction coefficient, F_n the normal contact load, k_d the contact stiffness, x the mass displacement, \dot{x} its velocity and z the internal slipping variable.

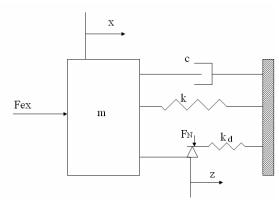


Figure 8: Masing's contact model in case of a 1 dof mass-spring system

Due to nonlinearity, solutions are obtained using a Newmark time integration scheme with Newton-Raphson iterations. Two different set of parameters (μ and k_d) are considered. The first is obtained using model updating from experimental data. The second considers estimated averaged values obtained from [2], as in this reference authors suggest values representing a best compromise when experimental data are not available: friction coefficient of 0.66 and contact stiffness of 3e7 N/m. Table 5 gives the parameters used for the smooth surface with an area of 13x13 mm² for models 1 and 2. Parameters suggested in [2] are higher than the identified values.

	Friction coefficient	Contact stiffness
Model 1	0.45	3.9e5 N/m
Model 2	0.66	3e7 N/m

Table 5: Parameter for the smooth 13x13 mm² surface

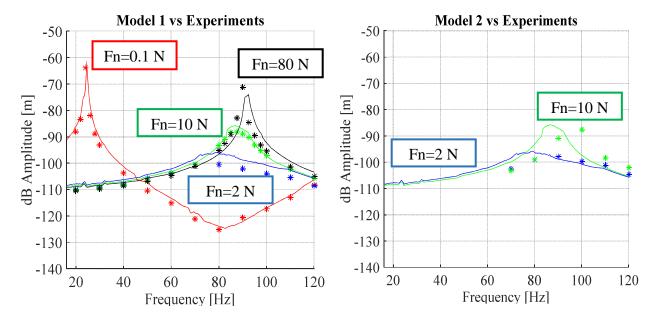


Figure 9: Models 1 and 2 versus experiments

Figure 9 gives numerical and experimental results. In the case of model 2, comparisons are limited to two normal loads, due to calculation time. Results clearly show that model 1 reproduces with good accuracy the experimental FRFs while there is a shift of 10 Hz between model 2 and experiments. Minimal discrepancy is obtained when reaching the maximal displacement amplitudes for both models. A 10 Hz frequency shift is not dramatic considering that no experimental data is available, but illustrates the need for accurate parameter estimation.

Numerical studies were made for the other configurations of contact area size and roughness and confirm little to no influence of these parameters on the dynamic response of the assembly. An example of comparison between surface 13x13 and 5x5 is shown in Table 6.

	Fn=0.1 N	Fn=2 N	Fn=10 N	Fn=80 N
Amplitude differences.	0 %	1.4 %	2.2 %	-3.1 %
Frequency of max amplitudes.	0 %	0 %	0 %	0 %
Damping ratio differences.		1.7 %	1.6 %	0.23 %

Table 6: Contact area size effect on the FRF – smooth 13x13mm² versus smooth 5x5 mm²

In order to examine a broader range of configurations, numerical calculations were performed for higher normal loads (120-160-200 N) and excitation loads (2-20-40-100 N). In Figure 10, smooth surface 13x13 and rough surface 13x13 are compared for Fn=120 N and Fex=40 N. The maximal amplitude discrepancy is 6.1 % and it appears that the smooth surface dissipates significantly more energy than the rough one. Other calculations show an increase of the influence of contact characteristics for higher external loads (Fn and Fex). However, the observed discrepancies for maximum amplitude remain limited and do not exceed 10% for instance.

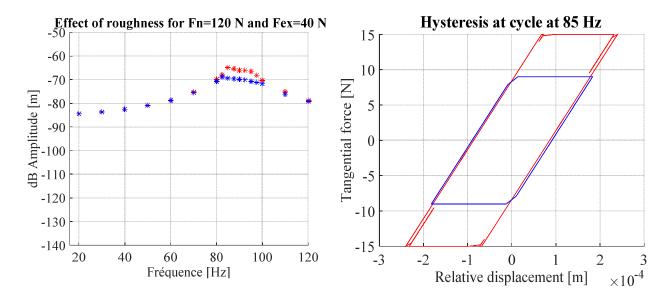


Figure 10: Smooth surfaces (red) versus rough surface (blue) for 13x13 sample, Fn=120 N and Fex=40 N.

5 CONCLUSIONS

This study focuses on the dynamics of a beam associated with rubbing devices. Both experimental and numerical results are presented. Different configurations were tested in order to assess the effect of contact area size and roughness of the assembly. It was shown that the influence of such characteristics is limited when considering the discrepancies on maximum amplitude of displacements, frequencies of peak displacements and damping ratios. For the configurations considered, the behaviour remain globally in the frame of macro-slip.

Numerically, calculations with Masing's model were done. In the case where experimental data is available, the numerical model and experiments agree very well and confirm the fact that contact area and roughness have limited effects on the FRFs. On the other hand, when measured data are not available, the correlation between the model (using averaged estimated values) and experiments is more delicate as a shift between the peak frequencies is observed. This is due to the fact that the contact stiffness estimation is too high. However a good agreement was noticed for maximum amplitudes. This last results emphasize that work for a better estimation of contact parameters is still needed. A model which does not require prior experimental friction tests such as in [5] is under development.

6 ACKNOWLEDGMENT

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